SUGRA without MSUGRA

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(work in progress)

Weak scale susy with a R parity is an elegant and cool solution to the hierarchy problem

But it is also a supercomplex beast

the complexity lies **not** in the number of additional particles introduced (superpartners + extra Higgs)

the complexity lies in the peripheral mechanisms that need to be introduced to make susy work

the root cause of all trouble:

SUSY needs to be broken

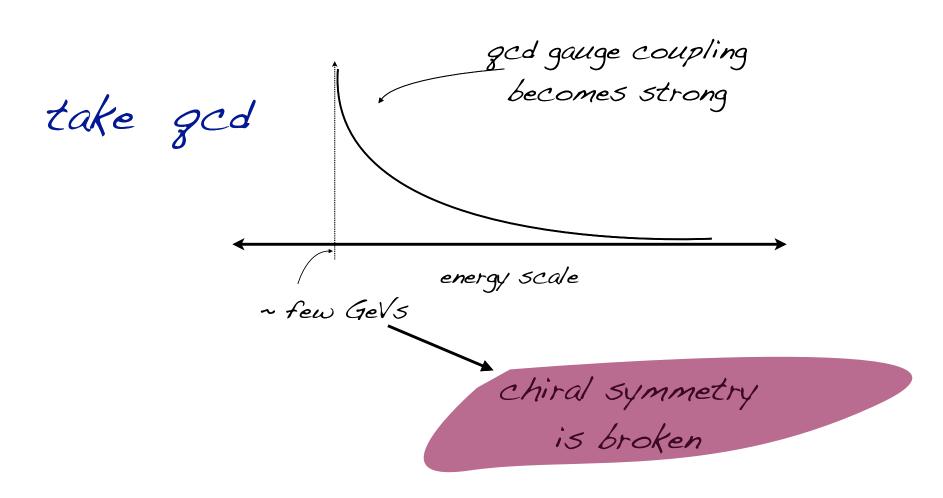
Trouble with breaking SUSY:

susy can't be broken by any of the SM fields or their superpartners because of the sum rule

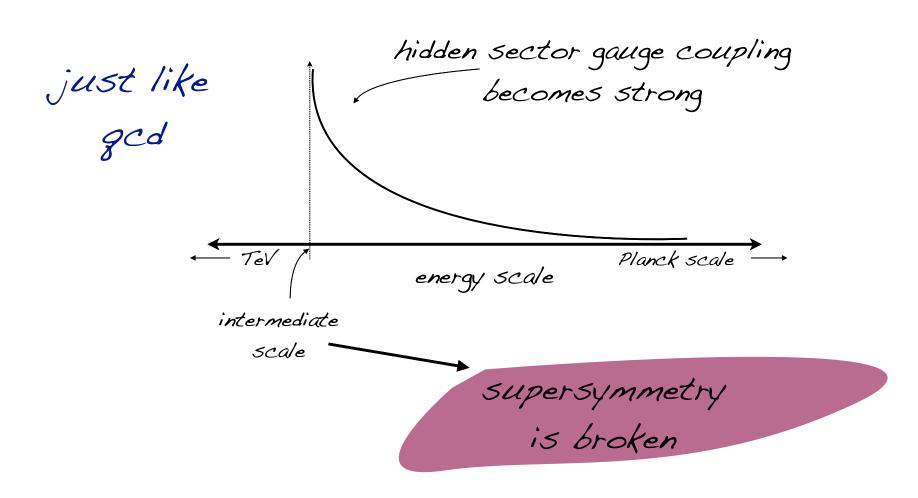
$$\operatorname{Tr}(\mathbf{m_s^2}) - 2\operatorname{Tr}(\mathbf{m_F^\dagger m_F}) + 3\operatorname{Tr}(\mathbf{m_V^2}) = 0$$

need to have a separate sector to break susy

SUSY is broken dynamically in a hidden sector



SUSY is broken dynamically in a hidden sector



Where is SUSY broken?

hidden sector

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How is SUSY broken?

DSB

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How is SUSY breaking communicated?

messenger mechanism gauge, gaugino, anomaly etc etc

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How is SUSY breaking communicated?

messenger mechanism gauge, gaugino, anomaly etc etc

 $\mu - B_{\mu}$ problem

tachyonic sleptons

Where is SUSY broken?

hidden sector

How is SUSY broken?

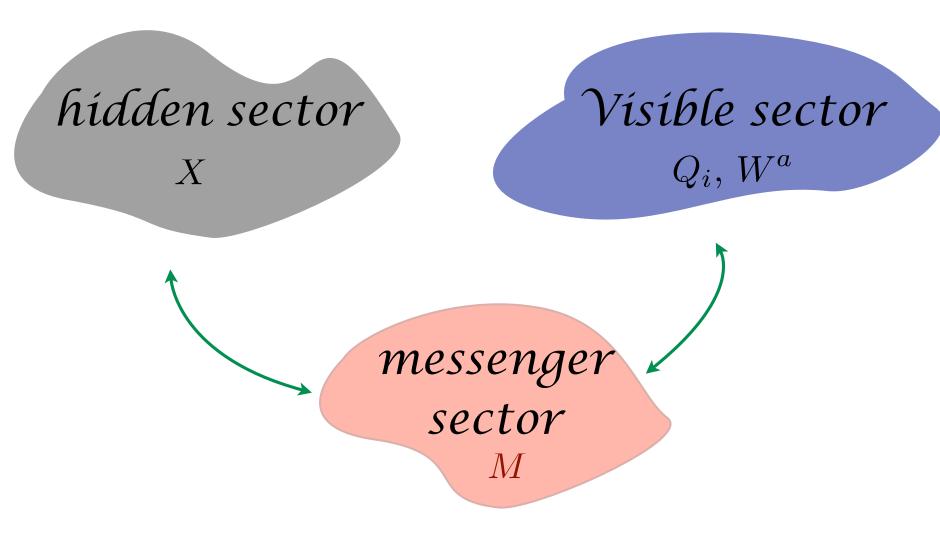
DSB

How is SUSY breaking communicated?

messenger mechanism gauge, gaugino, anomaly etc etc

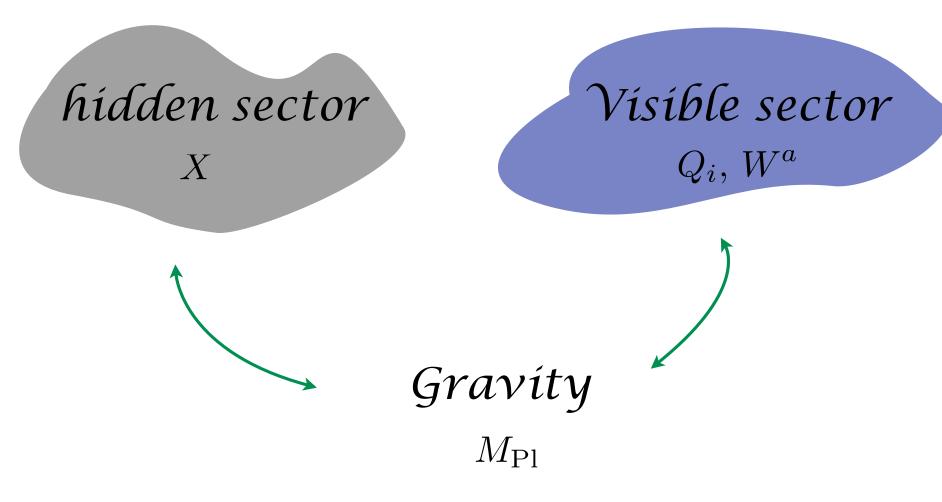
this is what will be simplified in this talk

general structure of a susy model



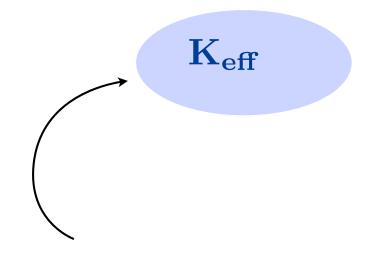
$$c_{ij} \frac{1}{M^2} \ X^{\dagger} X \ Q_i^{\dagger} Q_j$$

no need for additional messenger mechanism

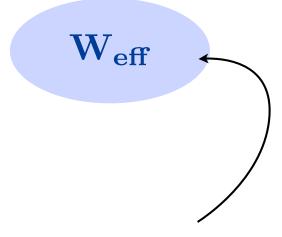


$$c_{ij} \frac{1}{M_{\rm Pl}^2} \ X^{\dagger} X \ Q_i^{\dagger} Q_j$$

after Planck scale physics is integrated out the effective interactions can be summarized in



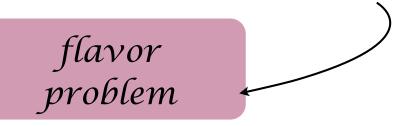
real and generic



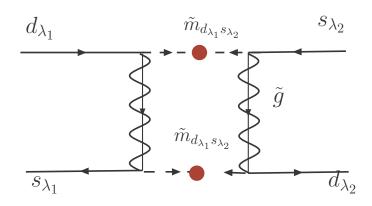
holomorphic and non-generic

after Planck scale physics is integrated out the effective interactions can be summarized in

real and generic



For example, $K_0 - \bar{K}_0$ mixing



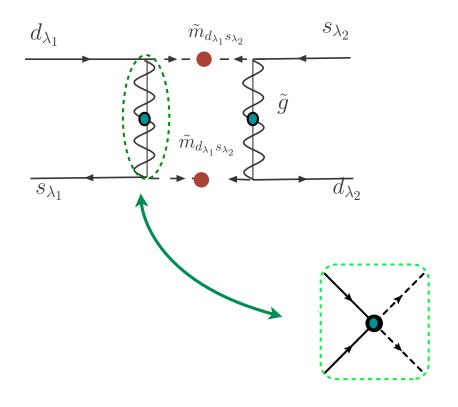
$$\delta_L \equiv \frac{(m_{\tilde{Q}}^2)_{12}}{m_{\tilde{q}}^2} \qquad \delta_R \equiv \frac{(m_{\tilde{D}}^2)_{12}}{m_{\tilde{q}}^2}$$

$$\delta_L = \delta_R = \delta \lesssim 0.003$$

$$\begin{cases} m_{\tilde{q}} = 1 \text{ TeV} \\ m_{\tilde{g}} = 2 \text{ TeV} \end{cases}$$

For example, $K_0 - \bar{K}_0$ mixing





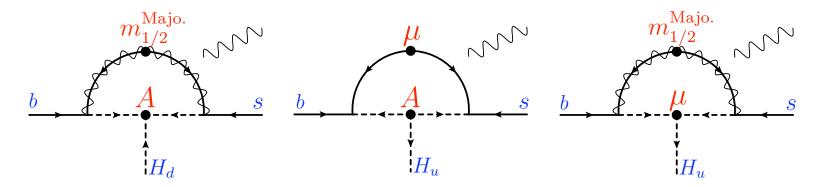
$$\frac{1}{m_{\tilde{g}}}(d_i d_j) \tilde{d}_k^* \tilde{d}_l^*$$

$$\Delta M_{K_0-ar{K}_0} \propto \left(rac{m_{ ilde{q}}}{m_{ ilde{q}}}
ight)^2$$

suppressed but not enough

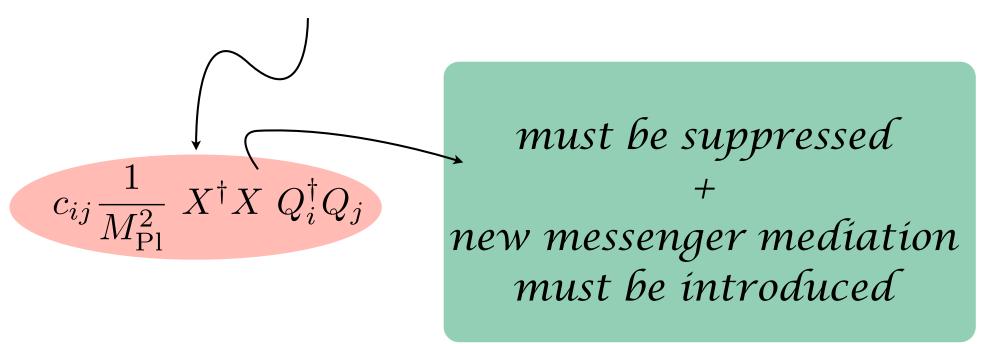
life is worse when we include Higgs sector

--- Excessive $\mu \to e \gamma$, $b \to s \gamma$ etc. ---



Flavor problem kills gravity mediation

but gravity mediation is always there



Here is our solution

old gravity mediation

step I
specific
hidden sector

step II tweak visible sector

Flow of the talk

current state of affair of susy model building

our solution

step I

step II

a sample point

phenomenology

conclusion

Step 1: hídden sector

only consider model which has no singlet and has an U(1) that participate in susy breaking

A concrete example:

(Dine, Nelson, Nir, Shirman, '95)

SU(4) X U(1)

$$X_1$$
 \square -3

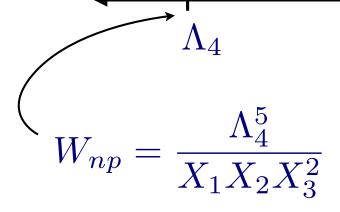
$$X_2$$
 \Box -1

$$X_3$$
 \square 2

$$X_4$$
 1 4

$$g_4 \gg \lambda \gg g_1$$

$$W = \lambda X_1 X_2 X_4$$



susy breaking vev at

$$D_1 = 1.4 \, F_{X_4} \neq 0$$

that was step I

but even before I proceed to step II let me briefly mention how life is significantly better ...

let's postpone Higgses for now

$$egin{aligned} \mathbf{K}_{\mathrm{eff}} &\supset c_{ij} rac{1}{M_{\mathrm{Pl}}^2} \ X^\dagger X Q_i^\dagger Q_j \end{aligned} \ egin{aligned} \mathbf{W}_{\mathrm{eff}} &\supset w_a rac{X}{M_{\mathrm{Pl}}} \ X_a W_a + \sqrt{2} \ \Omega_a rac{W'}{M_{\mathrm{Pl}}} W_a \Phi_a \end{aligned} \qquad egin{aligned} \mathcal{D}irac \\ mass \end{aligned}$$

let's postpone Higgses for now

$$\mathbf{K}_{\mathrm{eff}} \supset c_{ij} \frac{1}{M_{\mathrm{Pl}}^2} X^{\dagger} X Q_i^{\dagger} Q_j$$

$$\mathbf{W}_{\mathrm{eff}} \supset w_a \frac{X}{M_{\mathrm{Pl}}} \mathbf{W}_a W_a + \sqrt{2} \Omega_a \frac{W'}{M_{\mathrm{Pl}}} W_a \Phi_a$$

These interactions preserve an $U(1)_R$

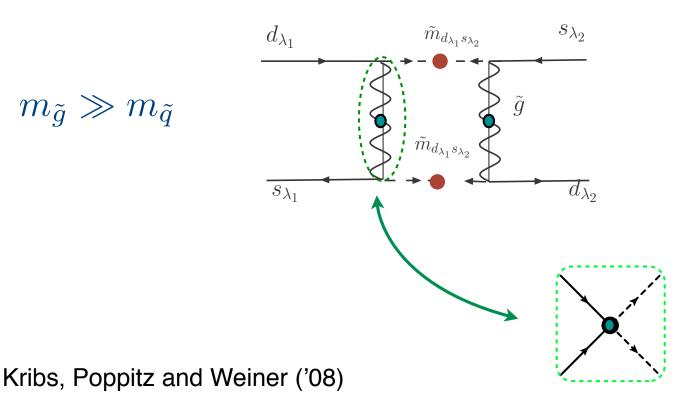
$$R[Q_i] = 1$$

$$R[W_a] = 1$$

$$R[\Phi_a] = 0$$

For example, $K_0 - K_0$ mixing



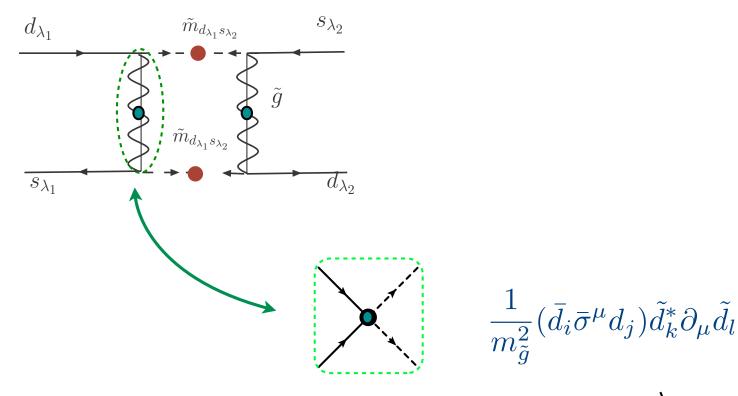


 $\frac{1}{m_{\tilde{g}}}(d_i d_j) \tilde{d}_k^* \tilde{d}_l^*$

No such operator would be generated

breaks R by 2 units

For example, $K_0 - \bar{K}_0$ mixing



enough suppression if

$$\frac{m_{\tilde{q}}}{m_{\tilde{g}}} \sim 0.1 \qquad \frac{\tilde{m}_{12}^2}{\tilde{m}_q^2} \sim 0.1$$

$$\Delta M_{K_0 - \bar{K}_0} \propto \left(\frac{m_{\tilde{q}}}{m_{\tilde{g}}}\right)^4$$

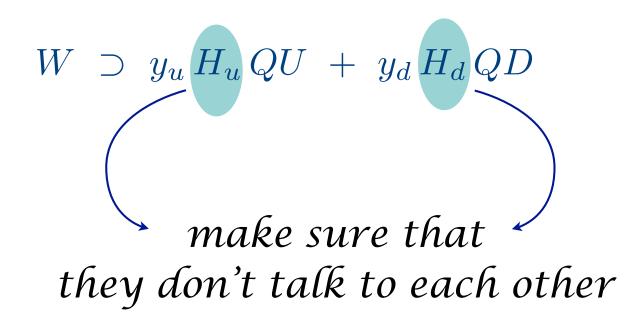
let's now include the Higgs sector

$$\mathbf{W}_{\mathbf{eff}} \supset \frac{1}{M_{\mathrm{Pl}}} X (Q \mathbf{W} H_u + Q \mathbf{N} H_d)$$
 $\mathbf{K}_{\mathbf{eff}} \supset \frac{1}{M_{\mathrm{Pl}}} X^{\dagger} \mathbf{W}_u H_d$

there is no A-term or μ term all flavor violation due to A & μ vanish

but Higgsinos are massless!

Prescription for Higgsino mass



Also, no singlet means no G-M mass term

step II

as promised I will now tweak the visible sector

step II visible sector

$$W \supset SHR + S^2T$$

EWSB potential is stable around

$$\langle S \rangle \neq 0, \quad \langle H \rangle \neq 0. \quad \langle R \rangle = \langle T \rangle = 0$$

no $X^{\dagger}X H^{\dagger}R$

new gauge

group

$$\mathbf{K}_{\text{eff}} \supset \frac{X^{\dagger}X}{M_{\text{Pl}}^{2}} \quad c_{ij}Q_{i}^{\dagger}Q_{j} + \\ \frac{X^{\dagger}X}{M_{\text{Pl}}^{2}} \quad (b_{h}H_{u}H_{d} + b_{r}R_{u}R_{d} + b_{s}S_{u}S_{d} + b_{t}T_{u}T_{d})$$

$$\mathbf{W}_{\mathrm{eff}} \supset \sqrt{2} \Omega_a \frac{W'}{M_{\mathrm{Pl}}} W_a \Phi_a$$

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$$\mathbf{W}_{\mathrm{eff}} \supset \sqrt{2} \Omega_a \frac{W'}{M_{\mathrm{Pl}}} W_a \Phi_a$$

$$R[H_{u,d}] = 0$$

$$R[R_{u,d}] = 2$$

$$R[S_{u,d}] = 0$$

$$R[T_{u,d}] = 2$$

$$\mathbf{K}_{\mathrm{eff}} \supset \frac{X^{\dagger}X}{M_{\mathrm{Pl}}^{2}} \quad c_{ij}Q_{i}^{\dagger}Q_{j} +$$

$$\frac{X^{\dagger}X}{M_{\rm Pl}^2} \quad (b_h H_u H_d + b_r R_u R_d + b_s S_u S_d + b_t T_u T_d)$$

$$\mathbf{W}_{\text{eff}} \supset \sqrt{2} \Omega_a \frac{W'}{M_{\text{Pl}}} W_a \Phi_a$$

breaks R by 4 unit

$$R[H_{u,d}] = 0$$

$$R[R_{u,d}] = 2$$

$$R[S_{u,d}] = 0$$

$$R[T_{u,d}] = 2$$

wouldn't give rise to dangerous dim-5 flavor violating operators since they break R by 2 unit

Flow of the talk

current state of affair of susy model building

our solution

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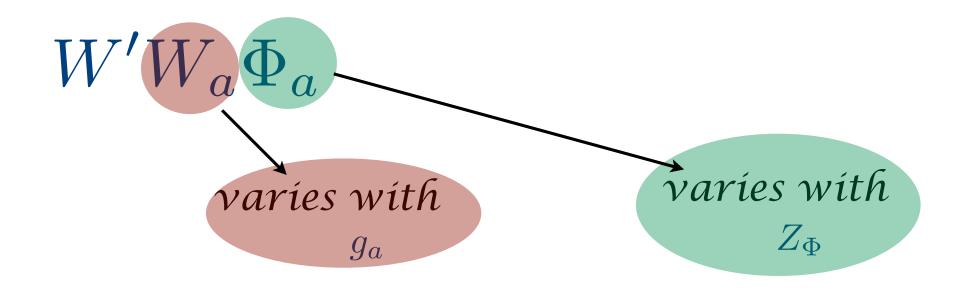
Planck scale inputs

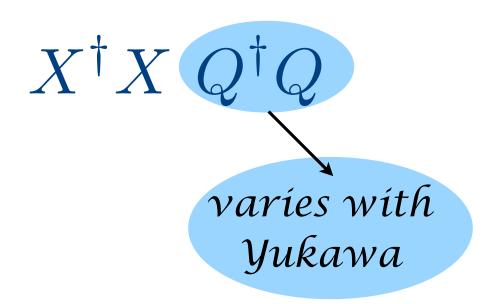
$$\frac{F_X}{M_{\rm Pl}} = \frac{1}{1.4} \frac{D}{M_{\rm Pl}} = 500 \text{ GeV}$$

$$\Omega_a = 1.4$$
 $M_a = 1 \text{ TeV}$ $c_{ij}^{\text{sq}} = 0.4$ $m_{\tilde{q}}^2 = (300 \text{ GeV})^2$

completely anarchic squark soft mass

RGE evolution



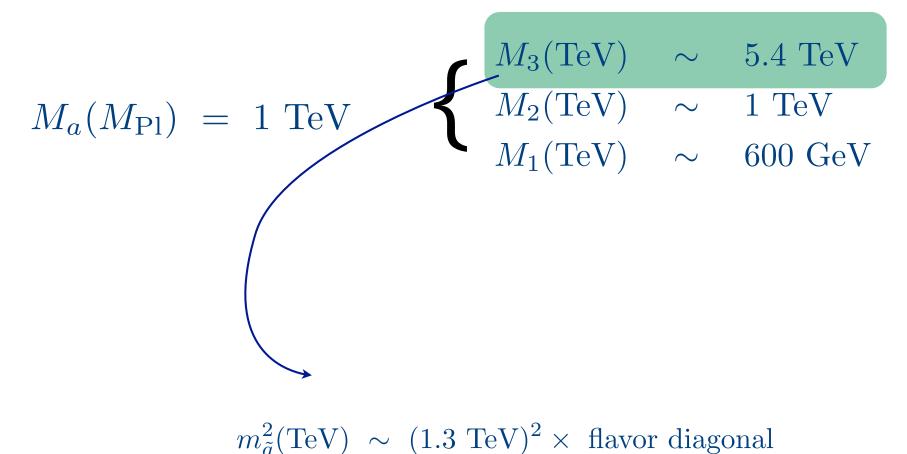


TeV scale outputs

$$M_a(M_{\rm Pl}) = 1 \,{
m TeV}$$

```
M_3({
m TeV}) \sim 5.4 {
m TeV}
M_2({
m TeV}) \sim 1 {
m TeV}
M_1({
m TeV}) \sim 600 {
m GeV}
```

TeV scale outputs



Friday, July 9, 2010

TeV scale outputs

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m GeV} \end{cases}$$

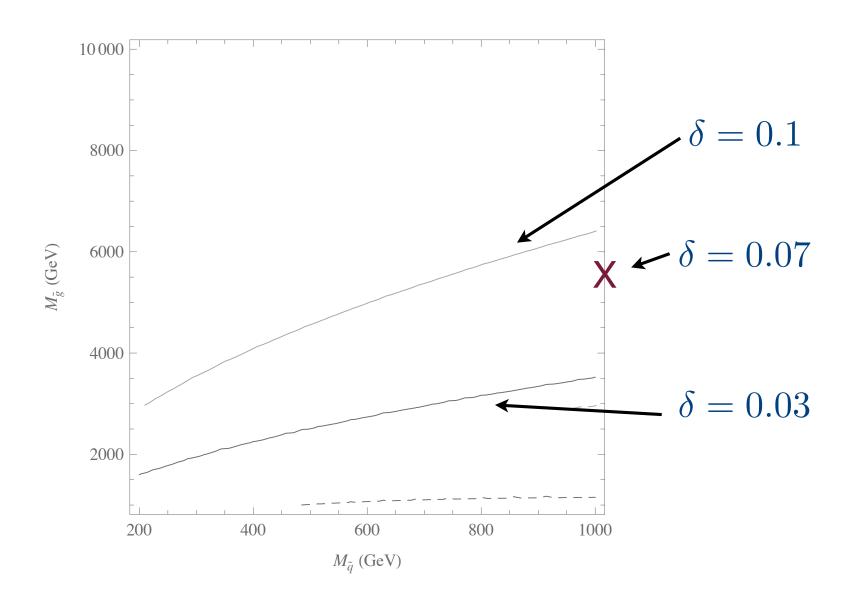
$$m_{\tilde{q}}^2(\text{TeV}) \sim (300 \text{ GeV})^2 \times \text{anarchic} + (1.3 \text{ TeV})^2 \times \text{flavor diagonal}$$

$$\frac{m_{\tilde{q}}}{m_{\tilde{g}}} \sim 0.2 \qquad \frac{\tilde{m}_{12}^2}{\tilde{m}_q^2} \sim 0.05$$

is it safe from $K_0 - \bar{K}_0$

Blechman and Ng ('08)

QCD corrections included



ϵ_K produces stronger bound

$$M_3({
m TeV}) \sim 5.4 {
m TeV}$$

 $M_2({
m TeV}) \sim 1 {
m TeV}$
 $M_1({
m TeV}) \sim 600 {
m GeV}$

$$m_{\tilde{q}}^2(\text{TeV}) \sim (300 \text{ GeV})^2 \times \text{anarchic} +$$

$$(1.3 \text{ TeV})^2 \times \text{flavor diagonal}$$

$$\theta \equiv \operatorname{Im}(\delta_L^* \delta_R) \lesssim .01$$

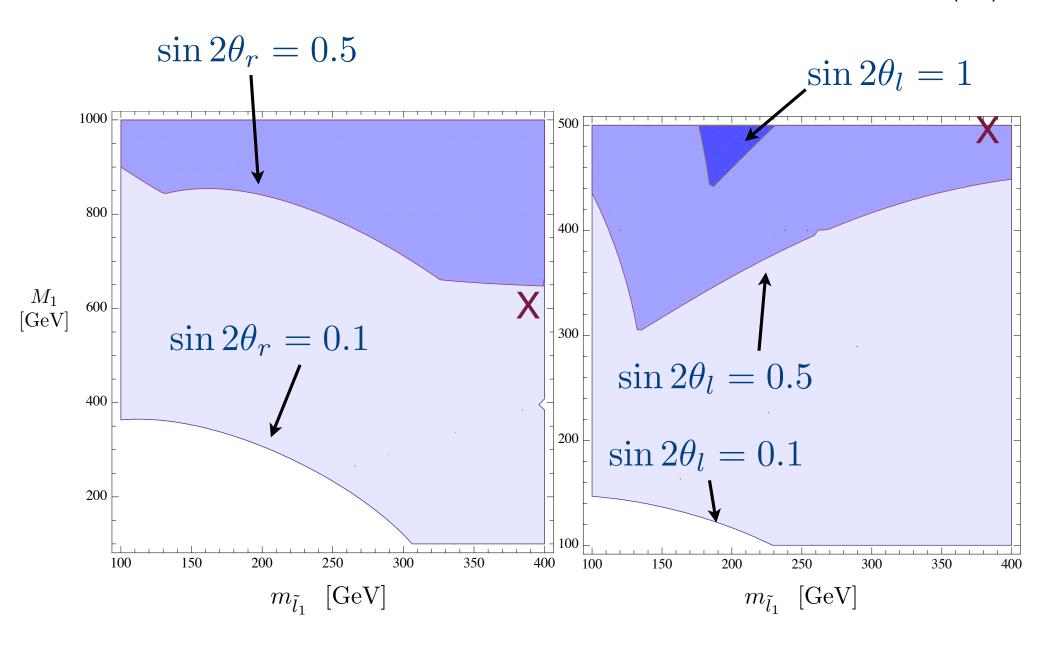
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 $M_2(\text{TeV}) \sim 1 \text{ TeV}$
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 $m_{\tilde{q}}^2(\text{TeV}) \sim (300 \text{ GeV})^2 \times \text{anarchic} + (1.3 \text{ TeV})^2 \times \text{flavor diagonal}$

$$\theta \equiv \operatorname{Im}(\delta_L^* \delta_R) \lesssim .01$$
0.2

lepton flavor constraints

Fok and Kribs ('10)

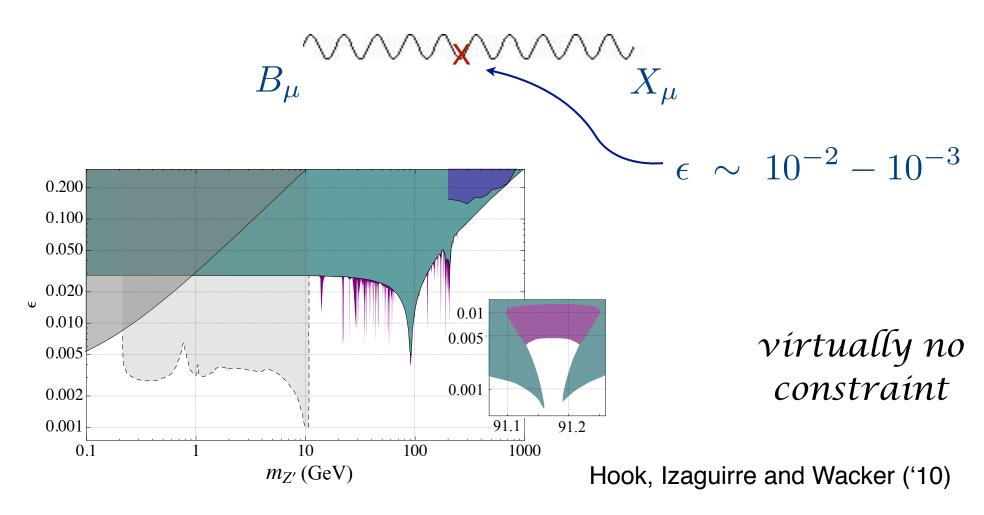


Phenomenology

- · gauginos are Dirac
- · flavor violation waiting to be discovered at LHC
- \cdot prospect of finding Z' at the EWSB scale

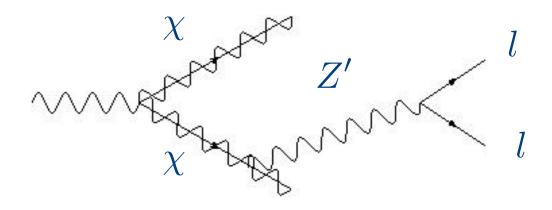
Z' at the $\mathcal{E}WSB$ scale

$M_{Z'}$ is constrained because of



Z'at the EWSB scale

Z' talks to Higgs sector with gx strength



Conclusions

- We resurrected gravity mediation
- Our prescription:
 - a singletless hidden sector with an U(1) that participates in susy breaking
 - Higgs sector is extended to include an $\mathrm{U}(1)_X$
- The effective theory contains an accidental (approximate) R-symmetry and RGE assisted hierarchy that solves the flavor problem
- ▶ Prospect of finding a Z' at the EWSB scale